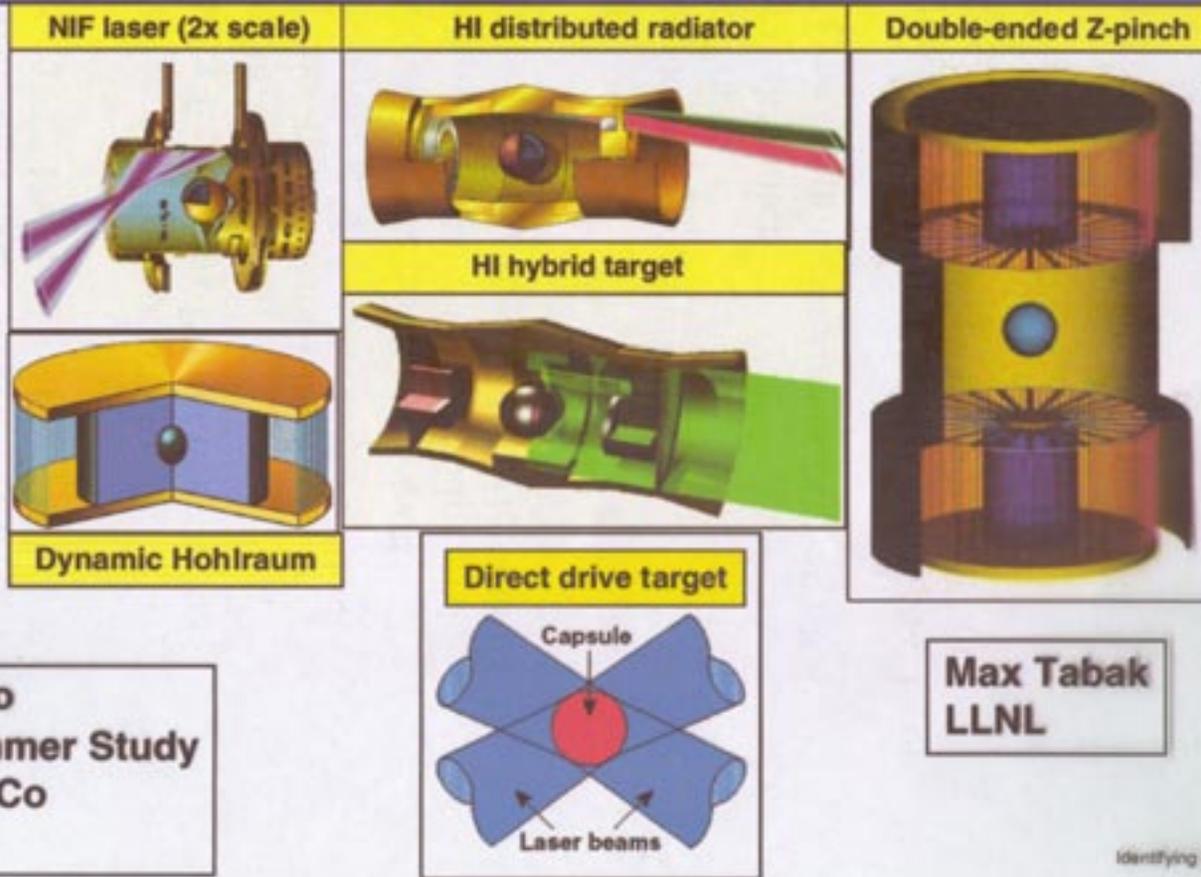


# Target Physics Issues for Inertial Fusion Energy

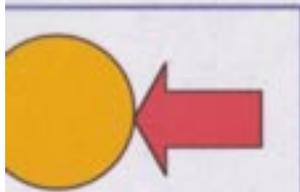


Presented to  
Fusion Summer Study  
Lawrence Livermore National Laboratory  
July 8, 2002

Identifying Mark-

## Target designs can be characterized by ignition method, compression method and driver

### Ignition

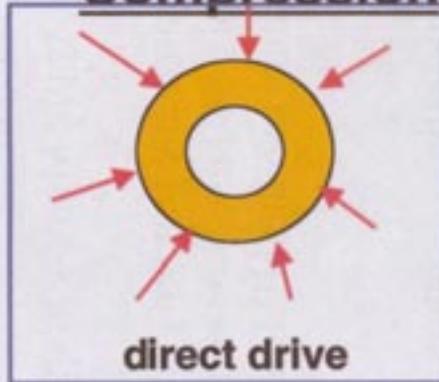


Fast Ignition

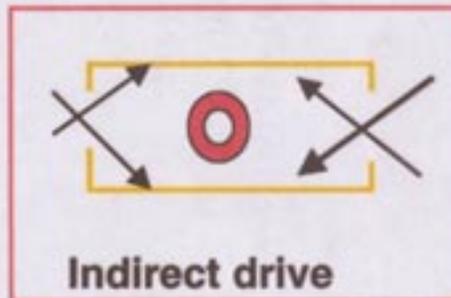


Ignition by  
stagnation of  
convergent flow

### Compression



direct drive



Indirect drive

### Driver

Laser  
 $\eta=5-10\%$

Heavy ion  
Accelerator  
 $\eta=15-40\%$

Z-pinch  
 $\eta\sim 15\%$

## How do targets interface with the rest of the IFE system?

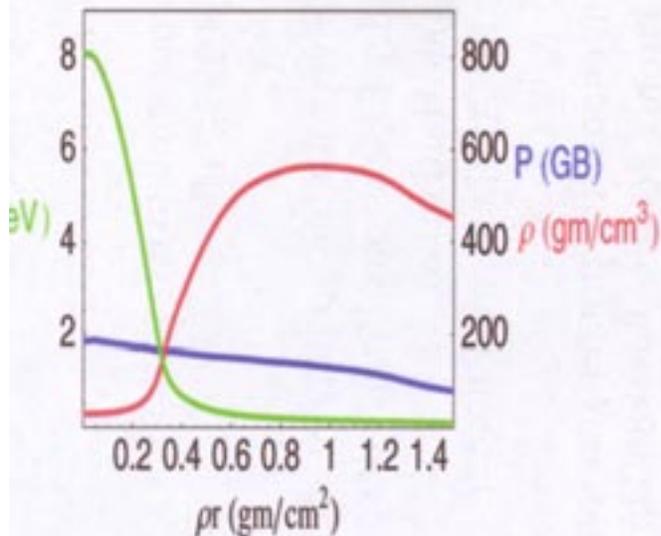
- Sets driver requirements
  - Scale and pulse shape (adiabatic control) required to meet G
  - Illumination geometry and precision: symmetry control
  - Beam quality and smoothness: hydro & plasma instabilities
  - Beam brightness: convergence ratio or Fast Ignition requirement
- Sets fabrication requirement
  - Ablator/DT smoothness : residual Rayleigh-Taylor instability
  - Hohlraum design & materials: symmetry control & efficiency
- Delivers photon, neutron and debris insult to first wall and optics

**Require sufficient gain(G) q.v.  $\eta G > 10$ : CoE**

## There are several factors that control gain

- **Optimal 1-D behavior**
  - Entropy management by pulse shaping and control of preheat
  - Minimum implosion velocity =  $v(\alpha, Pr, E)$
  - Rocket efficiency
- **Hydrodynamic instability growth must be minimized to control shell breakup and mixing of high(hotspot) and low(main fuel) entropy regions**
- **Implosion symmetry must be controlled so that all fuel elements stagnate simultaneously in a compact mass**
- **Coupling efficiency set by beam-plasma coupling and symmetry requirements**

The ignition energy can be estimated assuming an isobaric hot spot



Energy accounting:

$$E_c + E_h \sim \alpha_{stag}^{3/5} P_s^{2/5} M_c + C \frac{(\rho_h R_h)^3 T_h^3}{P_s^2}$$

optimizing over  $P_s$

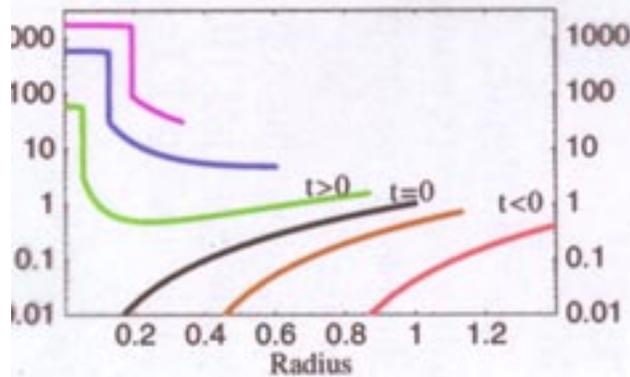
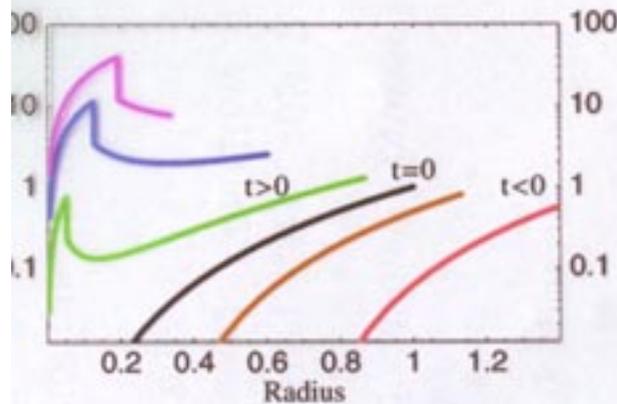
$$E_{ign} \sim \alpha_{stag}^{1/2} M_c^{5/6} \sim \alpha_{stag}^{1/2} \left( \frac{E_{ign}}{v^2} \right)^{5/6}$$

$$E_{ign} \sim \alpha_{stag}^3 v^{-10}$$

tailed numerical  
simulations give:

$$E_{ign} (kJ) = 2.1 \alpha_{stag}^{2.7} \left( \frac{v}{3 \cdot 10^7 \text{ cm/sec}} \right)^{-7.2}$$

# Self-similar implosions suggest the stagnation adiabat depends on $\alpha_{if}, v, P$



Meyer-ter-Vehn and Schalk found:

$$\rho_f \approx 2.4 \rho_0 M^{1.5} \quad P_f \approx 3.6 P_0 M^3$$

$$\alpha_{stag} \sim \frac{P_f}{\rho_f^{5/3}} \sim \alpha_{if} \sqrt{M}$$

$$\alpha_{stag} \sim \alpha_{if}^{0.85} v^{0.5} P^{-0.1}$$

From numerical simulations:

$$E_{ign} \text{ (kJ)} = 50.8 \alpha_{if}^{1.88} v^{-5.89} P^{-0.1}$$

$v$  in units of  $3 \cdot 10^7$  cm/sec

$P$  in units of 100MB

# The Rayleigh-Taylor instability is one of the key physics issue facing IFE

A target will fail if the perturbation amplitude (a) equals the shell thickness ( $\Delta R$ ).

Perturbations grow exponentially:

$$a_L(t) = a(t_0)e^{\gamma t}$$

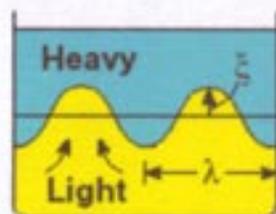
$$\gamma = \sqrt{Akg} - \beta k V_A$$

Ablative stabilization term

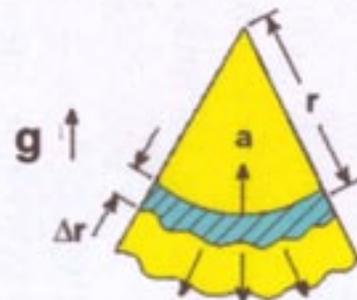
$$V_A = \frac{\dot{m}}{\rho_{\text{peak}}}$$

until saturation

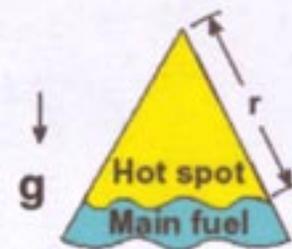
$$a(t) = a_{\text{sat}} \{1 + \log[a_L(t)/a_{\text{sat}}]\}$$



classical



acceleration



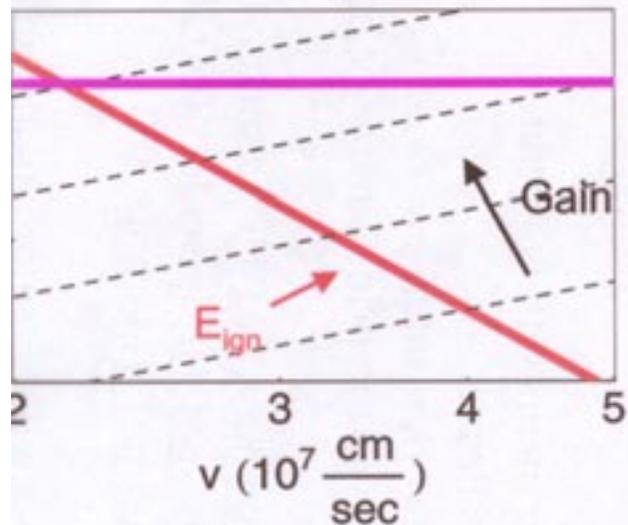
deceleration

# Capsule optimization involves trading off robustness versus perturbation growth



Time fixed adiabat, drive pressure, and fuel energy

Increasing Perturbation Growth  $\rightarrow$



For fixed capsule absorbed energy  
which velocity is best?

	Low Velocity	High Velocity
Perturbation Growth	Low	High
Gain	High	Low
Margin	Low	High

detailed numerical simulations with imposed surface roughness show the fastest shell is most susceptible to shell breakup



Slow



Moderate



Fast



Very Fast



$\rho$

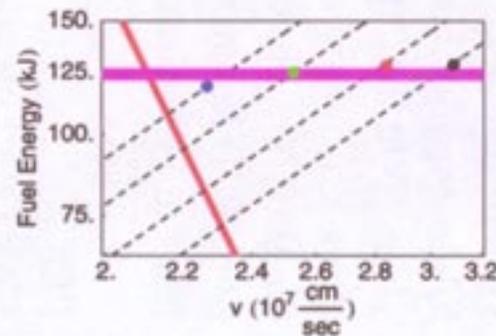


$$R_{\text{outer}} \sim 2.3 \text{ mm}$$

$$E_{\text{abs}} \sim 900 \text{ kJ}$$

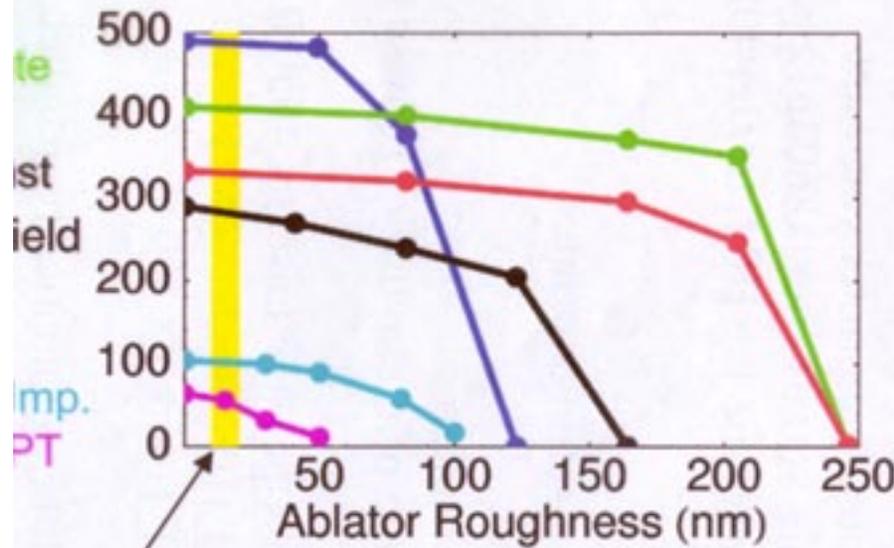
$$T_{\text{foot}} \sim 80 \text{ eV}$$

$$T_{\text{max}} \sim 265 \text{ eV}$$



All four capsules start with 8x NIF standard ablator roughness and standard ice roughness

# High-yield plastic capsules can tolerate ablator roughness ~10-20x NIF standard



NIF standard" surface finish

For this configuration:

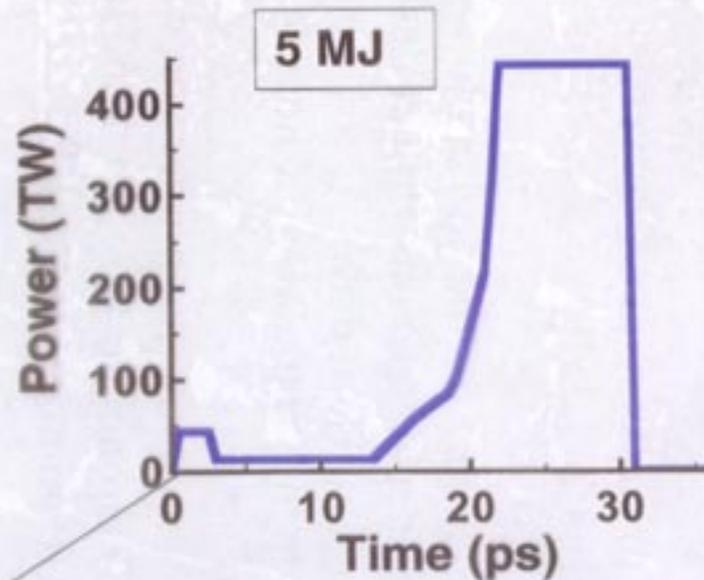
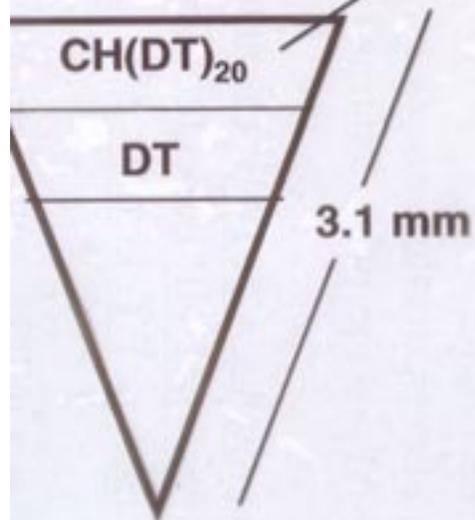
- Slow capsule fails due to cold fuel into the hot spot
- Very Fast capsule fails due to "shell breakup"
- Moderate and Fast capsules (Margin 1.3-1.6) are most robust
- Slow capsule gives highest gain for roughness < 80nm

**NIF and High Yield Capsules show similar tradeoffs between 1-D robustness and perturbation growth**

*pulse shaping*

High performance target designs combine wetted foam with adiabat shaping

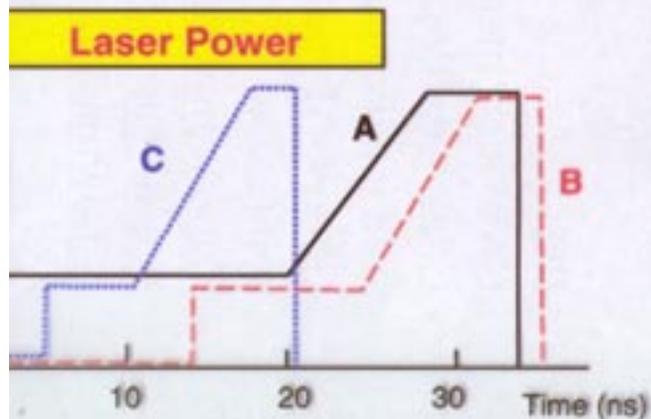
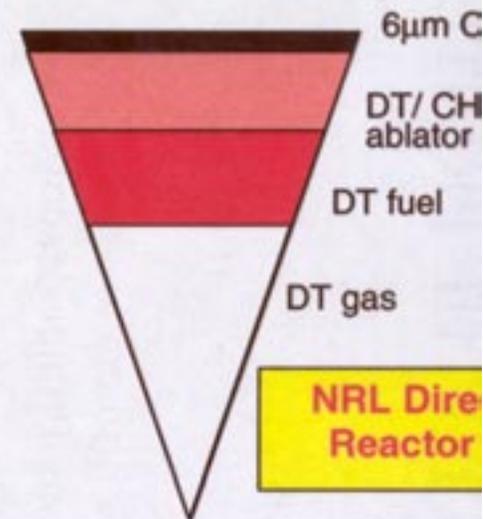
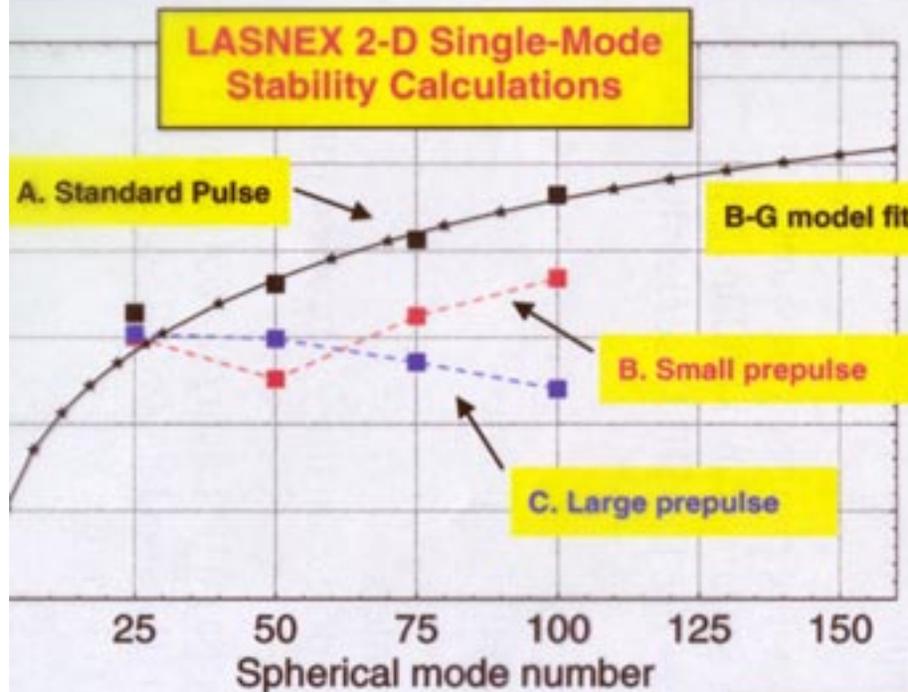
high-z material increases laser absorption



Intensity spike reduces imprint and shapes adiabat

Fig. 1

# Shaping the Laser Pulse Offers Large Improvements in Stability for the Same Direct Drive Target



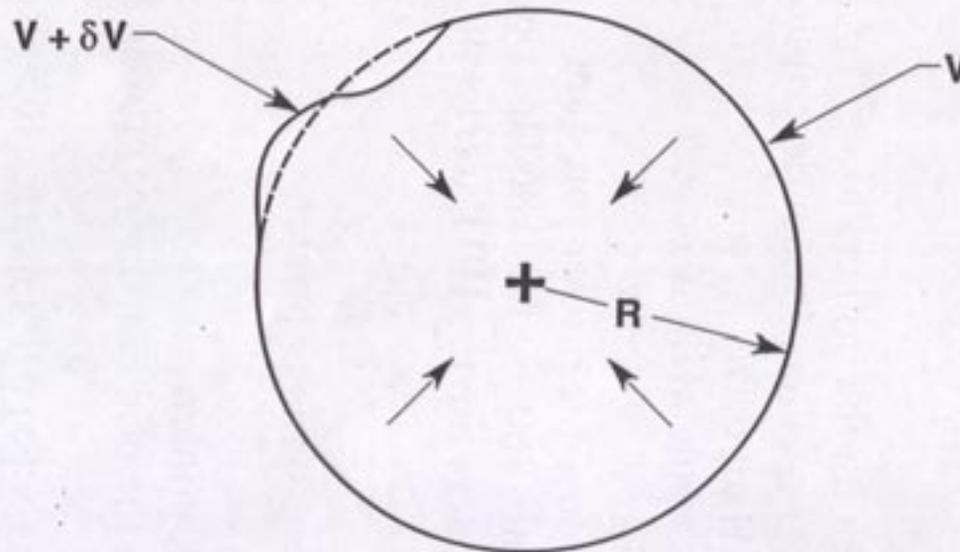
Pulse Shape	Laser (MJ)	Yield (MJ)	Gain	M Break
A. Standard	2.4	430	180	
B. Small prepulse (small "picket")	2.5	420	170	
C. Large prepulse (large "picket")	3.1	360	110	

With contributions from A. Schmitt (NRL), R. Be...

Symmetry

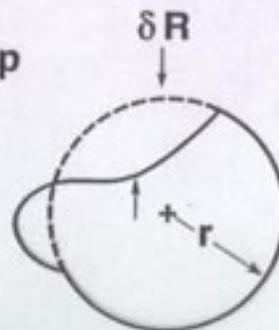
losion symmetry is an important issue for convergence-ratio targets

Small nonuniformity when outershell is at large radius



becomes magnified when shell is imploded to a very small radius

lower peak compression, temp  
lower  $\rho R$

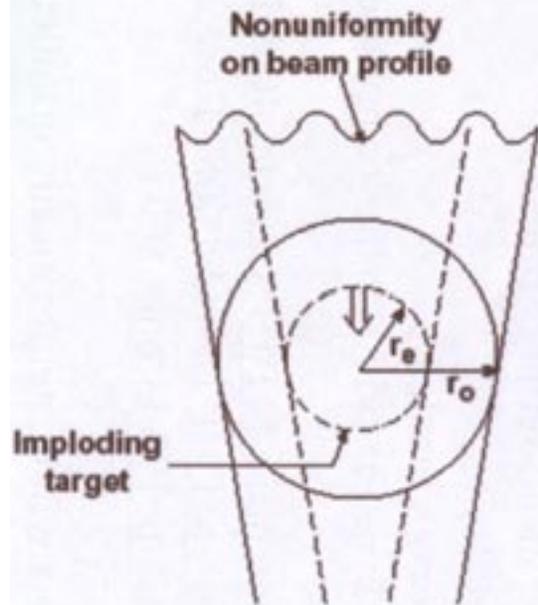
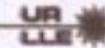


$$\delta R = (\delta V)t \sim \delta V \frac{R}{V} < 1/2 r$$

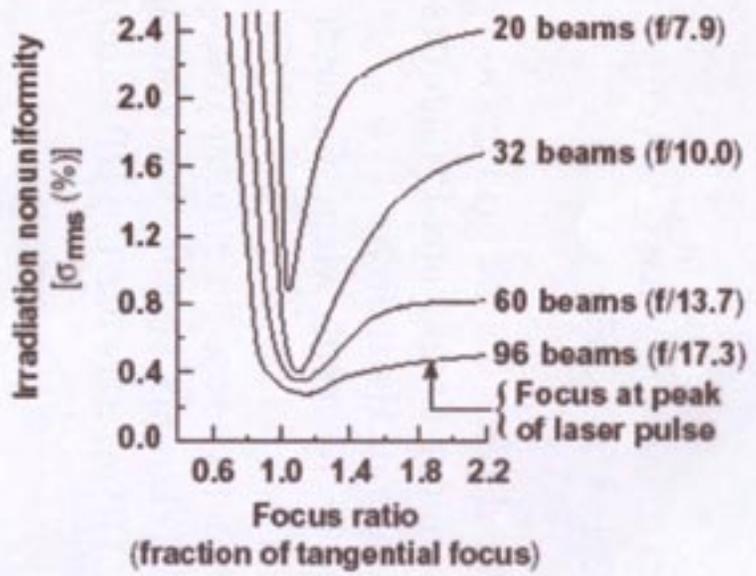
$$\therefore \frac{\delta R}{r} = \left( \frac{\delta V}{V} \right) \frac{R}{r} < 1/2$$

$$\therefore \frac{\delta V}{V} < 1/2 \frac{r}{R} < 1/2 \text{ (conv. ratio)}$$

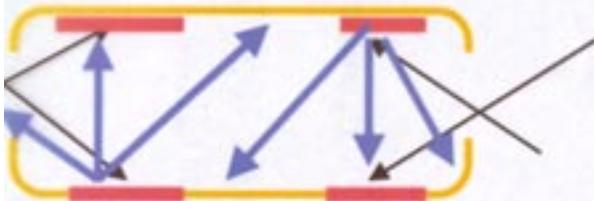
# A sufficient number of beams provide adequate long-wavelength uniformity throughout the implosion



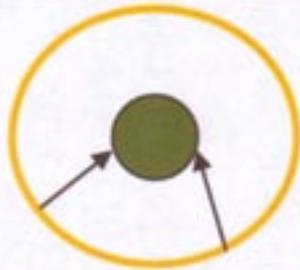
$r_o$  – initial radius  
 $r_e$  – radius at end of laser pulse



## Radiation symmetry in indirectly driven targets is controlled by three factors



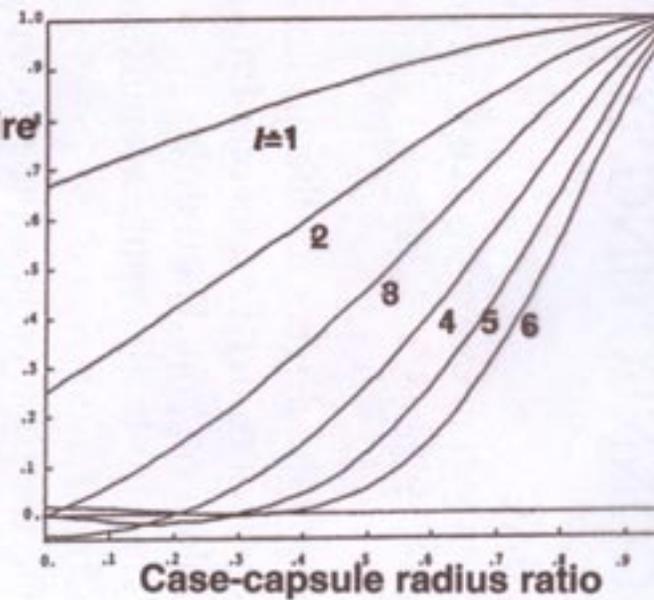
Hohlraum re-radiation smooths wall temperature by factor  $1/(1-\alpha)$



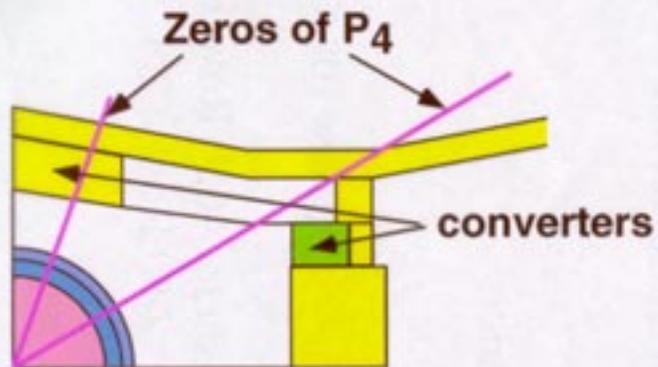
Ratio of Legendre amplitudes

to ball transport smooths  $l$  modes

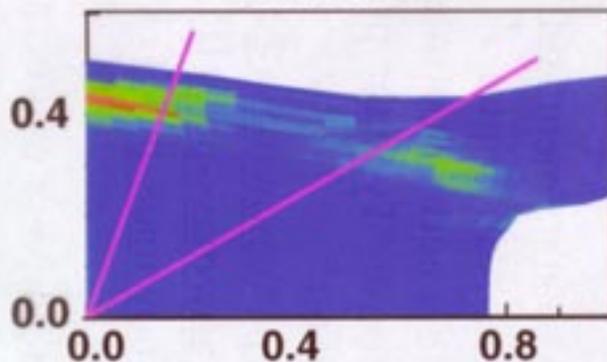
Appropriate location of beam spots  
suppresses low order modes and allows  
large case-capsule area ratios



For the heavy ion distributed radiator target, symmetry is achieved by depositing most of the energy near the zeros of  $P_4$



Deposited energy density at 28 nsec

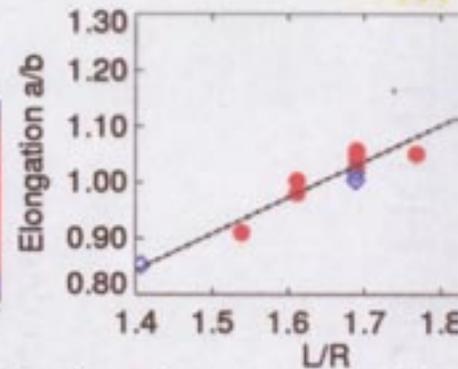
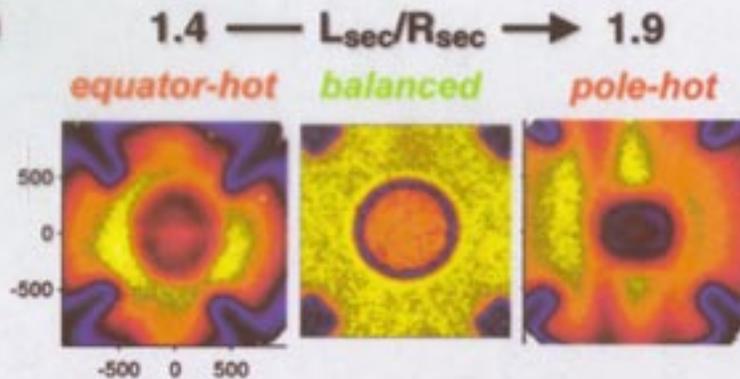


- For a spherical hohlraum, sources with Gaussian weights located at the zeros of  $P_4$  causes all moments from  $P_1$  to  $P_7$  to be zero
- Converters are placed near the zeros of  $P_4$  so that most of the energy is deposited there
- Converter densities are adjusted so that  $P_2$  is small

# Demonstrating capsule symmetry control is a high priority for near-term experiments on Z

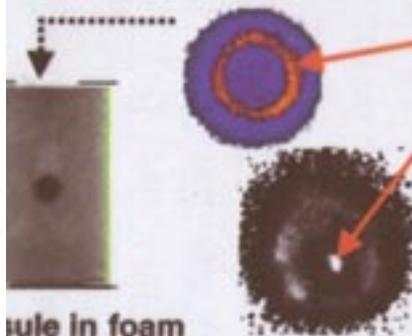


## Double-Pinch Hohlraum



- Symmetry control via hohlraum design (e.g.  $L_{sec}$  variation)
- Agreement with viewfactor-RHD + integrated calculations
- Optimum length symmetry is close to scaled HY requirements

## Dynamic Hohlraum

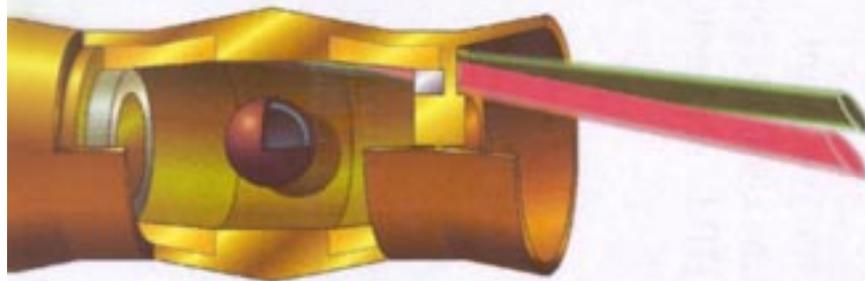


- Shock emission width consistent with low pinch  $R_{sec}$
- End-on core x-ray images + spectra, side-on planar
- Foam dopants, radiation burnthrough shielding, and converter shaping for pole-equator asymmetry

capsule in foam

© 2002 WRA/Vesey V1

**Using a smaller hohlraum allows us to trade  
coupling efficiency and symmetry**



**Coupling efficiency: 17%**  
**Gain: 66**  
**Driver Energy: 5.9 MJ**

**Hohlraum dimensions reduced by  
6% but driving the same capsule**

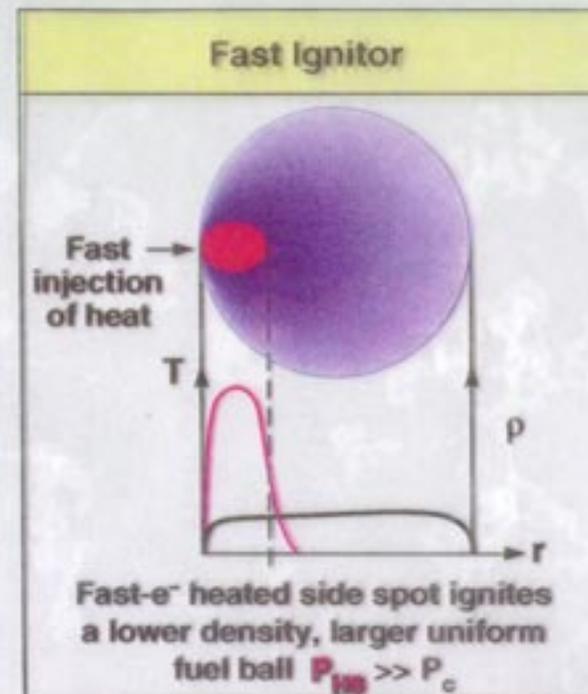
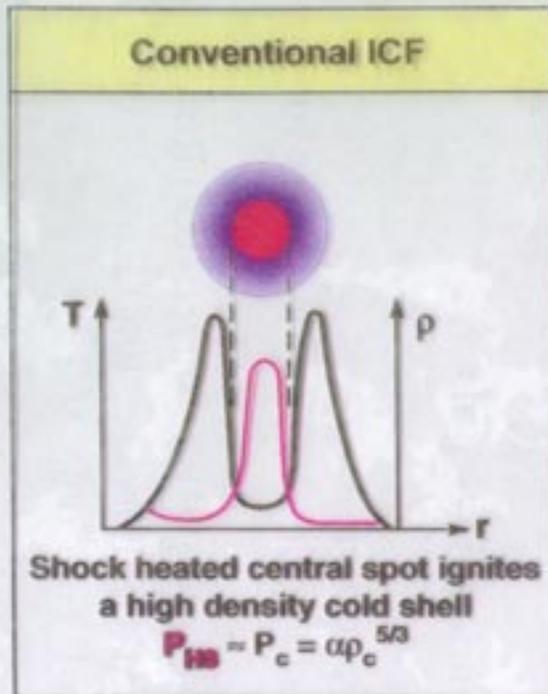


**Coupling efficiency: 27%**  
**Gain: 130**  
**Driver Energy: 3.3 MJ**

## **Beam plasma coupling seems acceptable but studies continue**

- **Laser coupling efficiency 70%(direct drive)-90%(indirect drive or zooming)**
  - **Study SBS in hohlraums**
- **Filamentation can change location of laser energy deposition,affecting symmetry and driving other instabilities**
- **Raman scattering can lead to hot,preheating electrons but levels seem to be acceptable**
  
- **Ion beam stopping power has 20% theoretical uncertainties**
  - **Plasma target experiments beginning**
  - **By varying target designs we can accommodate factor 2-3 uncertainty**

Technology advances had made innovative concepts possible: ultra-high brightness lasers may allow a fundamentally new method of igniting inertial fusion capsules



\* Tabak, Hammer, Glinsky, Krueer, Wilks, Woodworth, Campbell, & Perry *Phys. Plasmas* 1, 1626 (1994).

\*\* H. Azechi et al., *Laser Part. Beams* 9, 2 (1991).

### Advantages of Fast Ignitor

- Fast Ignitor implosions are less stressing: (mix, convergence, ...)
- Lower  $\rho \Rightarrow$  more mass to burn ( $E_c = \alpha M_c \rho_c^{2/3}$ )  $\Rightarrow$  Higher Gain

Significant R&D is required to explore potential of this concept

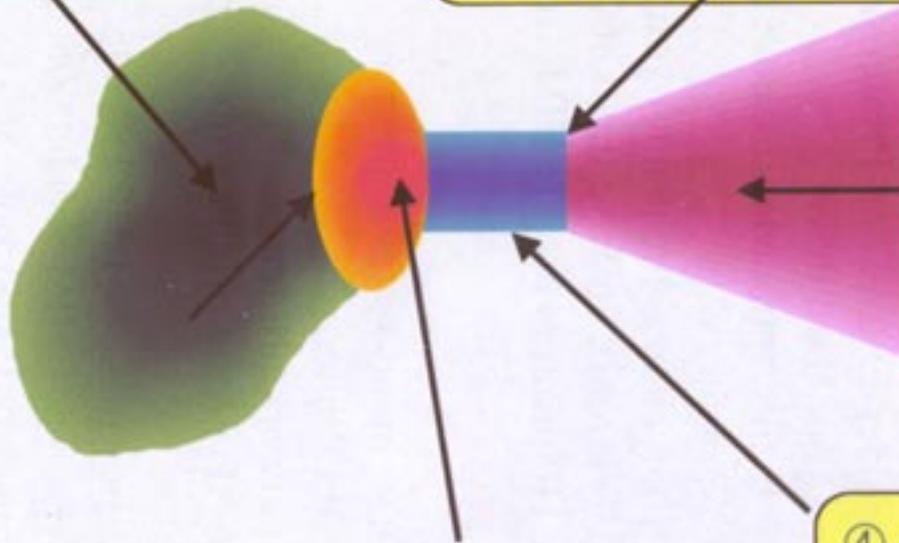
## Ignition requires several key elements:

① Implosion to high density blob

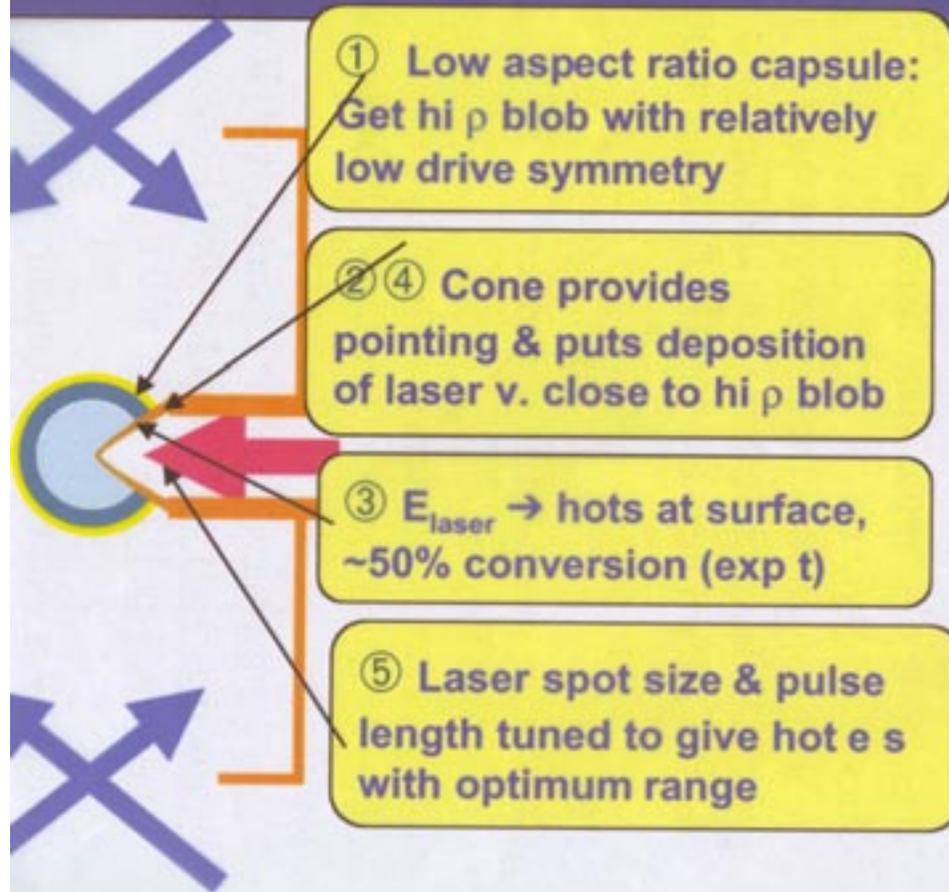
③ Efficient conversion:  
 $E_{\text{laser}} \rightarrow$  beam of  $e^-$ s or  $p^+$ s with appropriate range aimed at blob

② Short pulse laser aimed and timed to blob implosion

④ Efficient transport of hot electron(or proton) energy to blob



-focussed scheme provides  
t of the key FI elements.



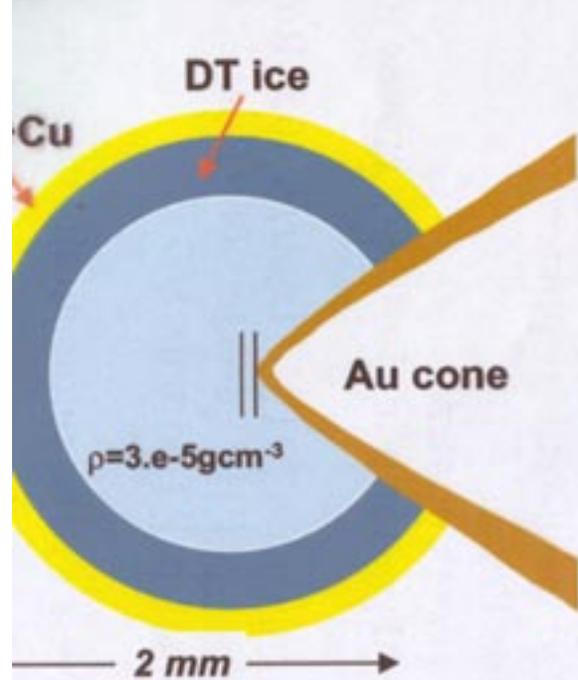
**Other options:**

**Direct drive**

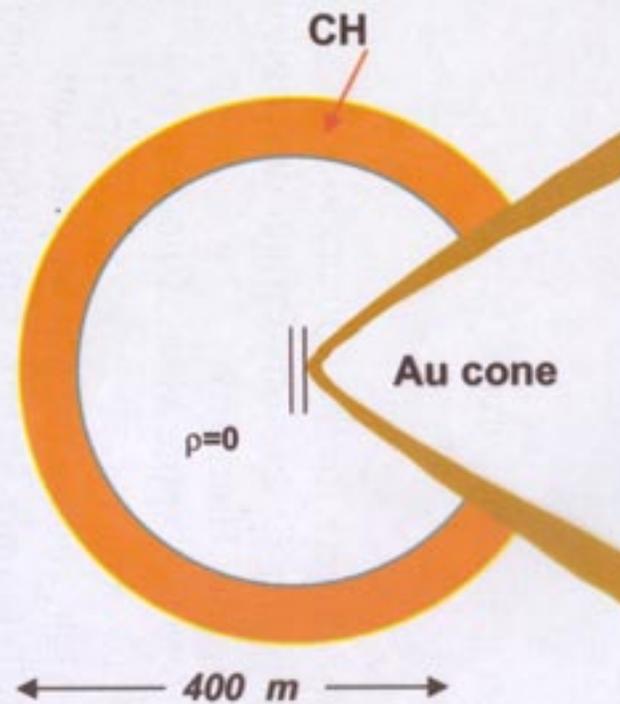
Use produced  
proton beam as  
intermediary

Without cone:  
Ponderomotive  
hole boring through  
coronal plasma

Omega experiments tested cone-focus concept for  
fast ignition at 1/5 scale

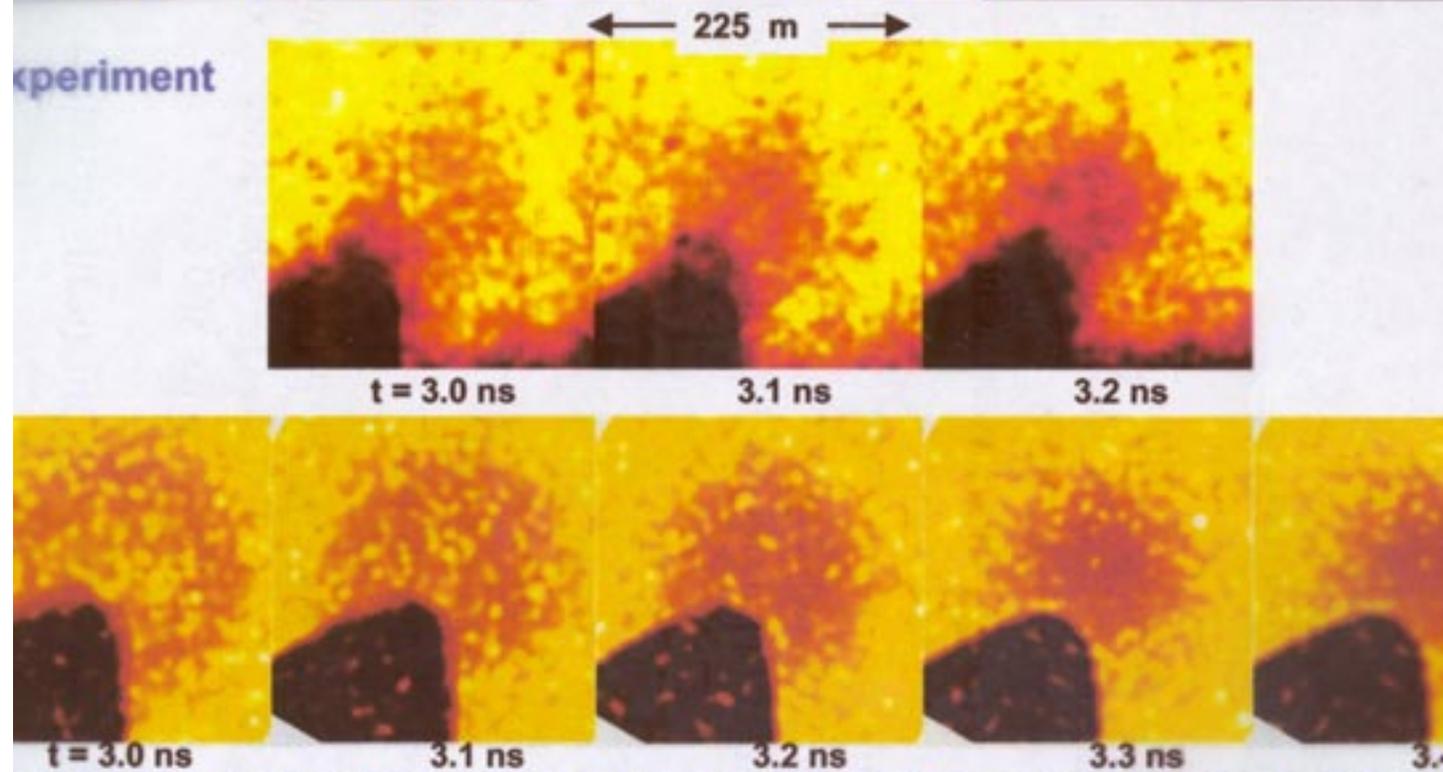


Ignition scale



Omega scale

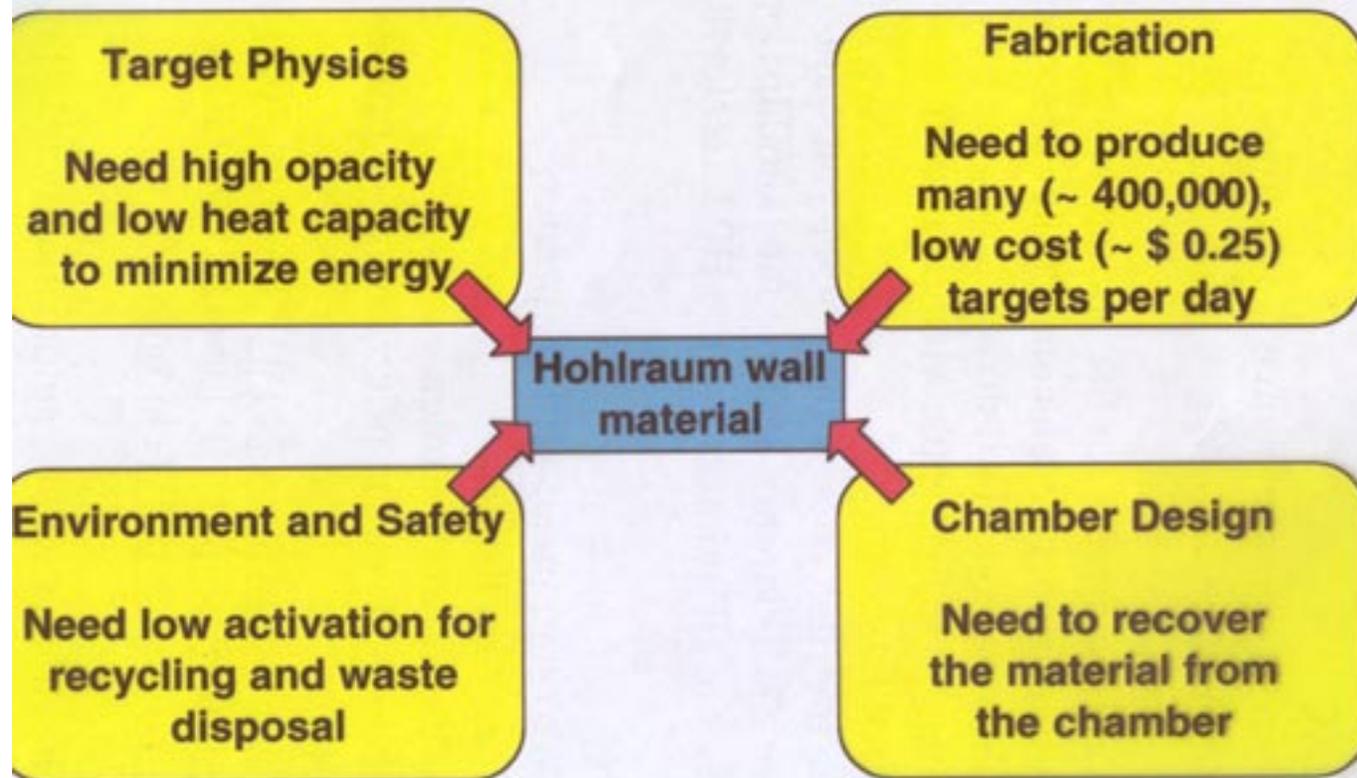
Backlit images (@8 keV) show convergence of one-focussed targets was very similar to prediction — with perhaps a small time offset.



prediction (with pixelation, noise, and smoothing like exp. images)

Comparison shows some exp. evidence for gold entrainment near tip of

## The hohlraum wall material must satisfy constraints from four areas



target designs with varying degrees of risk  
 provide adequate gain for all driver concepts

